

HFS Atomic Scattering Factors*

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Atomic scattering factors calculated from carefully computed Hartree–Fock–Slater wave functions are presented for the neutral elements from $Z=2$ to $Z=100$ over a range of $\sin \theta/\lambda$ values up to 6.0. Because of the nature of the wave functions used and because of the range and density of the parameters for which the calculations have been made, this compilation should represent the most complete and most accurate f values available.

Introduction

The need for good values of the atomic scattering factors in crystallographic studies has long been appreciated. With the introduction of new refined measurement techniques and with the efforts to solve more difficult and more subtle crystal structure problems, this need is now being felt even more acutely. It is generally conceded that the best values of the atomic scattering factors are those obtained from the most sophisticated and presumably most realistic atomic models. In particular, the Hartree–Fock (HF) values have found particular favor although the difficulty of obtaining true HF wave functions has precluded the compilation of form factors for anything but a limited number of elements. A simplification of the basic HF problem which was introduced by Slater (1951) provides one with a computationally tractable problem. The resulting Hartree–Fock–Slater (HFS) wave functions are very good approximations to those obtained by HF; whatever differences exist lie mainly in the same direction as the effect produced by further physical refinements of the HF problem.

Self consistent wave functions have been calculated by Herman & Skillman (1962) for all ground state neutral atoms from $Z=2$ to $Z=103$ using the HFS approach. These values represent the most complete atomic wave functions available both from the point of view of numerical accuracy and of atomic number. For this reason it was felt that a calculation of atomic scattering factors from these wave functions would be of interest and importance.

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Discussion

The calculations of the atomic scattering factors were made from the HFS wave functions by incorporating the assumption of spherical symmetry into the expression,

$$f = \sum_j \int |\psi_j|^2 \exp(i\kappa \mathbf{s} \cdot \mathbf{r}_j) d\tau_j, \quad (1)$$

obtaining

$$f = \sum_j \int_0^\infty U_j(r) \frac{\sin sr}{sr} dr \quad (2)$$

using a common notation.

The integration was carried out in a straightforward fashion by means of Simpson's rule. The r coordinate is expressed in terms of the Thomas–Fermi parameter,

$$r = \left(\frac{3}{32\pi^2} \right)^{\frac{1}{3}} \frac{\hbar^2}{2me^2} \frac{1}{Z^{\frac{1}{3}}} x = \frac{0.88534138}{Z^{\frac{1}{3}}} x$$

and the integration is performed over x when it has been expressed in units of the first Bohr radius. The x mesh is common to all atoms and is that utilized in the original HFS wave function calculations. It is constructed as follows: The first 41 values of x from 0.00 to 0.10 are given a constant spacing of 0.0025. The next 41 values of x range from an x of 0.10 to an x of 0.30 inclusive with a constant spacing of 0.005. The process of doubling is continued from one block of 41 mesh points until 441 values of x are involved. After expressing all distances in terms of the Bohr radius, the integration was performed with a suitable modification of Simpson's rule to account for the non-constant intervals. To check the accuracy of the wave

Table 1. Atomic scattering factors for elements Z=1 to Z=100 as a function of $\sin \theta/\lambda$

Z	Elements 1-10										Elements 11-20										Elements 21-30										Elements 31-40										Elements 41-50										Elements 51-60										Elements 61-70										Elements 71-80										Elements 81-90										Elements 91-100																																																																																																																																																																																													
	H	He	Li	Be	B	C	N	O	F	Ne	Na	Mg	Al	Si	P	S	Cl	Ar	K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr	Rb	Sr	Y	Zr	Mo	Tc	Ru	Rh	Pd	Au	Hg	Pb	Bi	Po	At	Rn	Fr	Ra	Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Mn	Co	Ni	Cu	Zn																																																																																																																																																																																																																	
.00	1.000	2.000	3.000	4.000	5.000	6.000	7.000	8.000	9.000	10.000	11.000	12.000	13.000	14.000	15.000	16.000	17.000	18.000	19.000	20.000	.00	21.00	22.00	23.00	24.00	25.00	26.00	27.00	28.00	29.00	30.00	31.00	32.00	33.00	34.00	35.00	36.00	37.00	38.00	39.00	40.00	.00	21.00	22.00	23.00	24.00	25.00	26.00	27.00	28.00	29.00	30.00	31.00	32.00	33.00	34.00	35.00	36.00	37.00	38.00	39.00	40.00																																																																																																																																																																																																																										
.05	.947	1.855	2.722	3.512	4.277	5.027	5.757	6.478	7.185	7.875	8.548	9.204	9.844	10.468	11.077	11.671	12.250	12.814	13.364	13.900	.05	20.20	21.24	22.27	23.28	24.23	25.16	26.07	26.96	27.84	28.71	29.57	30.42	31.26	32.09	32.91	33.72	34.52	35.30	36.07	36.83	37.58	38.32	39.05	39.77	40.48	41.18	41.87	42.54	43.20	43.85	44.49	45.12	45.74	46.35	46.95	47.54	48.12	48.69	49.25	49.80	50.34	50.87	51.39	51.90	52.40	52.89	53.37	53.84	54.30	54.75	55.20	55.64	56.07	56.50	56.92	57.33	57.74	58.14	58.53	58.92	59.29	59.65	60.01	60.36	60.70	61.03	61.35	61.66	61.96	62.25	62.53	62.81	63.08	63.34	63.59	63.84	64.08	64.31	64.53	64.75	64.96	65.16	65.35	65.53	65.71	65.88	66.04	66.19	66.34	66.48	66.61	66.74	66.86	66.98	67.09	67.20	67.30	67.40	67.49	67.57	67.65	67.72	67.79	67.85	67.91	67.96	68.01	68.06	68.11	68.16	68.20	68.24	68.28	68.32	68.36	68.39	68.42	68.45	68.48	68.51	68.54	68.56	68.58	68.60	68.62	68.64	68.66	68.68	68.69	68.71	68.72	68.73	68.74	68.75	68.76	68.77	68.78	68.79	68.80	68.81	68.82	68.83	68.84	68.85	68.86	68.87	68.88	68.89	68.90	68.91	68.92	68.93	68.94	68.95	68.96	68.97	68.98	68.99	69.00	69.01	69.02	69.03	69.04	69.05	69.06	69.07	69.08	69.09	69.10	69.11	69.12	69.13	69.14	69.15	69.16	69.17	69.18	69.19	69.20	69.21	69.22	69.23	69.24	69.25	69.26	69.27	69.28	69.29	69.30	69.31	69.32	69.33	69.34	69.35	69.36	69.37	69.38	69.39	69.40	69.41	69.42	69.43	69.44	69.45	69.46	69.47	69.48	69.49	69.50	69.51	69.52	69.53	69.54	69.55	69.56	69.57	69.58	69.59	69.60	69.61	69.62	69.63	69.64	69.65	69.66	69.67	69.68	69.69	69.70	69.71	69.72	69.73	69.74	69.75	69.76	69.77	69.78	69.79	69.80	69.81	69.82	69.83	69.84	69.85	69.86	69.87	69.88	69.89	69.90	69.91	69.92	69.93	69.94	69.95	69.96	69.97	69.98	69.99	70.00

Table 1 (cont.)

Table with 20 columns representing elements (Cb, Mo, Tc, Ru, Rh, Pd, Ag, Cd, In, Sn, Pm, Sm, Eu, Gd, Tb, Dy, Ho, Er, Tm, Yb) and rows of scattering factor data for various elements and isotopes.

isolated atoms. Thus empirically the HFS f values may be at least as good as the HF in the majority of the experimentally important cases. There are other inherent limitations and defects in these calculations. For example: (1) they do not make allowances for the dispersion effects which occur when λ is in the vicinity of an X-ray absorption edge, (2) they assume spherical symmetry, (3) the wave functions from which the f values are computed are not corrected for spin-orbit effects, and (4) the wave functions are not corrected for relativistic effects.

Conclusions and comments

Despite the intrinsic approximations in these calculations, it is felt that the scattering factors presented here represent a substantial improvement over most values available. They possess the inherent advantages of being relatively complete and of having been obtained in a coherent fashion, the same type of approximation being involved in the calculations for each of the atoms. There are numerous obvious ex-

tensions of these calculations. In preparation at present are the following: (1) the scattering amplitudes for electron scattering both within the first Born approximation and in the case where consideration is taken of the phase change, (2) the inelastic scattering factors, (3) X-ray and electron scattering factors for ions, and (4) analytic approximations to the X-ray and electron scattering factors.

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Hartree Scattering Factors for Elements 37 through 98*

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Hartree scattering factors are given for atoms of atomic number 37 through 98 and for several of their ions. The values are listed in tabular form and as analytic functions. Values are also given for two low-lying excited states of cerium. Differences between the Hartree and Thomas-Fermi-Dirac scattering factors are discussed.

Introduction

Atomic scattering factors obtained from Hartree self-consistent-field (SCF) calculations have been available in the past for only a very few atoms and ions above atomic number 36. Boyd, Larson & Waber (1963) have recently completed Hartree SCF calculations for all atoms of the periodic table, many ions, and many excited states. In this paper we present scattering factors for the ground states of elements of atomic number 37 and higher, for several of their ions, and for two excited states of cerium. (For elements of lower atomic number, scattering factors which have been obtained from more accurate wave functions are already available in standard references.)

We have published in a Los Alamos Scientific Laboratory report (Cromer, Larson & Waber, 1963) a set of Hartree SCF scattering factors for all atoms in the periodic table, in a more detailed form than in the present paper. Although this publication duplicates light element work by others, it has the advantage of presenting in a single unit complete information on Hartree SCF scattering factors.

Calculations

Details of the Hartree SCF calculations will be given elsewhere by Boyd, Larson & Waber (1964). The scattering factors were computed from the total radial density functions by straightforward means (James, 1948). Calculations were performed for the most part with an IBM 7090 computer and, in the

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